

1 **The Impact of Large-Scale Orography on Northern Hemisphere Winter**  
2 **Synoptic Temperature Variability**

3 Nicholas J. Lutsko\*

4 *Department of Earth, Atmospheric, and Planetary Sciences, Massachusetts Institute of*  
5 *Technology, Cambridge, Massachusetts*

6 Jane W. Baldwin

7 *Princeton Environmental Institute, Princeton University, Princeton, New Jersey*

8 Timothy W. Cronin

9 *Department of Earth, Atmospheric, and Planetary Sciences, Massachusetts Institute of*  
10 *Technology, Cambridge, Massachusetts*

11 *\*Corresponding author address: Nicholas Lutsko, Department of Earth, Atmospheric, and Plane-*  
12 *tary Sciences, Massachusetts Institute of Technology, Cambridge, Massachusetts*

13 E-mail: lutsko@mit.edu

## ABSTRACT

14 The impact of large-scale orography on wintertime near-surface (850hPa)  
15 temperature variability on daily and synoptic time-scales (days to weeks) in  
16 the Northern Hemisphere is investigated. Using a combination of theory,  
17 idealized modeling work and simulations with a comprehensive climate  
18 model, it is shown that large-scale orography reduces upstream temperature  
19 gradients, in turn reducing upstream temperature variability, and enhances  
20 downstream temperature gradients, enhancing downstream temperature vari-  
21 ability. Hence the presence of the Rockies on the western edge of the North  
22 American continent increases temperature gradients over North America  
23 and, consequently, increases North American temperature variability. By  
24 contrast, the presence of the Tibetan Plateau and the Himalayas on the  
25 eastern edge of the Eurasian continent damps temperature variability over  
26 most of Eurasia. However, Tibet and the Himalayas also interfere with the  
27 downstream development of storms in the North Pacific storm track, and thus  
28 damp temperature variability over North America, by approximately as much  
29 as the Rockies enhance it.

30  
31 Large-scale orography is also shown to impact the skewness of down-  
32 stream temperature distributions, as temperatures to the north of the enhanced  
33 temperature gradients are more positively skewed while temperatures to the  
34 south are more negatively skewed. This effect is most clearly seen in the  
35 northwest Pacific, off the east coast of Japan.

## 36 **1. Introduction**

37 Temperature variability is one of the most important features of the climate for human society  
38 and natural ecosystems, affecting, among many other things, agricultural and economic produc-  
39 tion (Lazo et al. (2011); Wheeler and von Braun (2013); Shi et al. (2015); Jahn (2015); Bathi-  
40 any et al. (2018)) and the rhythms of ecological seasons (Jackson et al. (2009); Bowers et al.  
41 (2016)). Changes in temperature variability may be among the most impactful aspects of future  
42 climate change, which has motivated much recent work on the mechanisms controlling temper-  
43 ature variability in present and future climates, with two primary foci: (1) the question of how  
44 Arctic amplification will influence mid-latitude temperature variability; and (2) the question of  
45 what controls the zonal-mean variance and higher-order moments of the temperature distribution  
46 (e.g., Schneider et al. (2015); Garfinkel and Harnik (2017); Linz et al. (2018)). With respect to  
47 Arctic amplification, it is now clear that, in winter, mid-latitude zonal-mean temperature variance  
48 will be reduced (Screen (2014); Schneider et al. (2015); Hoskins and Woollings (2015)), though  
49 the effect of Arctic amplification on higher moments of mid-latitude temperature distributions is  
50 still uncertain (e.g., Cohen et al. (2014); Barnes and Polvani (2015)).

51 Little work, however, has been done to understand what controls regional (zonally-asymmetric)  
52 patterns of temperature variability, despite their societal relevance. For instance, changes in heat-  
53 waves with global warming can be well predicted by superposing a mean shift on present-day  
54 daily temperature variability, so that understanding the pattern of temperature variance is key for  
55 forecasting spatial variations in heat-wave changes with warming (Rahmstorf and Coumou (2011);  
56 Lau and Nath (2012); Lau and Nath (2014); Huybers et al. (2014); McKinnon et al. (2016)).

57 An example of a regional difference in temperature variability can be seen in panels a and b  
58 of Figure 1. Whether using daily data (panel a) or filtering to synoptic time-scales (days to

59 weeks, panel b), North America experiences substantially more near-surface (850hPa) tempera-  
60 ture variability than Eurasia during boreal winter (December-January-February, DJF, see section  
61 2 for description of observational dataset). This is also shown by Figure 1e, which plots a longi-  
62 tudinal profile of DJF synoptic temperature variance at 50°N: temperature variance at this latitude  
63 is roughly twice as large over North America as over Eurasia. Investigating the contribution of  
64 large-scale Northern Hemisphere orography (Asian orography, which includes the Himalayas, the  
65 Tibetan Plateau and the Mongolian Plateau, and the Rockies) to the enhancement of temperature  
66 variability over North America compared to Eurasia is the primary goal of the present study.

67 Our analysis is based on the dominant control of winter synoptic temperature variability by hor-  
68 izontal advection, which implies in turn that mean horizontal temperature gradients, particularly  
69 meridional gradients, are the primary control on synoptic temperature variability (Schneider et al.  
70 (2015); Holmes et al. (2016); see section 3a below). It can be seen in panels c and d of Figure 1 that  
71 both zonal and meridional temperature gradients are larger over North America than over Eurasia  
72 during winter, suggesting that whatever causes these enhanced gradients is also responsible for the  
73 enhanced variability over North America. Specifically, the importance of temperature gradients  
74 for synoptic temperature variability implies a close link between the Northern Hemisphere winter  
75 stationary wave pattern and the regional distribution of winter temperature variability.

76 Waves forced by large-scale orography are a key component of the winter stationary wave pattern  
77 in the Northern Hemisphere (Held et al. 2002). Below, we show that orography increases down-  
78 stream temperature gradients and decreases upstream temperature gradients, with corresponding  
79 impacts on temperature variability. We demonstrate this mechanism in simulations with two ideal-  
80 ized atmospheric general circulation models (GCMs), one dry and one moist, which also allow us  
81 to investigate how the shape of the orography influences its impact on temperature variability and  
82 how moist processes impact the dynamics (section 3). We then present simulations with a com-



83 prehensive climate model in which the major Northern Hemisphere mountain ranges are flattened,  
84 to quantify the impact these have on winter temperature variability (section 4). A complicating  
85 factor is orography’s effect on downstream development: the presence of large-scale orography  
86 can weaken downstream eddies by interfering with the recycling of energy from upstream, leading  
87 to reduced temperature variability far from the orography.

88 By enhancing and reducing mean temperature gradients, orography also impacts the skewness  
89 of temperature distributions, which we explore in section 5. We end with conclusions in section 6.

## 90 **2. Data and Methods**

### 91 *a. Observational Data*

92 Observational data are taken from the Modern-Era Retrospective Analysis for Research and  
93 Applications (MERRA) dataset (Rienecker et al. 2011). The MERRA grid has  $1.25^\circ$  resolution  
94 in latitude and longitude, and we have taken daily-averaged data from December, January and  
95 February for the years 1979 to 2012.

### 96 *b. Dry GCM*

97 The dry GCM is the GFDL spectral dynamical core, which solves the primitive equations for  
98 a dry ideal gas on the sphere, and is forced by Newtonian relaxation to a prescribed zonally-  
99 symmetric equilibrium temperature field and damped by Rayleigh friction near the surface. The  
100 parameter settings are the standard Held-Suarez parameters with forcing symmetric about the  
101 equator (Held and Suarez 1994). This set-up produces an equinoctial climate similar to that of the  
102 real atmosphere, though there are no stratospheric polar vortices due to the uniform stratospheric  
103 relaxation temperature.

104 As in Lutsko and Held (2016), the model is perturbed by adding a Gaussian mountain, with the  
105 form

$$h(\phi, \lambda) = H \exp \left\{ - \left[ \frac{(\phi - \phi_0)^2}{\alpha^2} + \frac{(\lambda - \lambda_0)^2}{\beta^2} \right] \right\}, \quad (1)$$

106 where  $H$  is the maximum height of the mountain in meters;  $\lambda$  and  $\phi$  are longitude and latitude,  
107 respectively;  $\lambda_0$  and  $\phi_0$  are the co-ordinates of the center of the mountain; and  $\alpha$  and  $\beta$  are half-  
108 widths, both set to  $15^\circ$  in the main suite of simulations.  $\lambda_0$  and  $\phi_0$  were set to  $90^\circ\text{E}$  and  $45^\circ\text{N}$ ,  
109 respectively, in all simulations.

110  $H$  was varied from 333m, which is in the “linear” regime, with air mostly flowing up and over  
111 the mountain, to 4km, which is in the “non-linear” regime, with air mostly deflected around that  
112 orography (Lutsko and Held 2016). In every simulation, the model was run at T85 resolution with  
113 30 evenly spaced sigma levels, and the instantaneous wind, surface pressure and temperature fields  
114 were sampled once per day. We present results from simulations lasting 5000 days, with data taken  
115 from the final 4000 days.

### 116 *c. Moist GCM*

117 The moist GCM is the gray-radiation model first described by Frierson et al. (2006), though  
118 we have used the parameter settings of O’Gorman and Schneider (2008), and also included their  
119 parameterization of short-wave absorption by the atmosphere. The model uses the GFDL spectral  
120 dynamical core, and includes the simplified Betts-Miller (SBM) convection scheme of Frierson  
121 (2007). We show results using a convective relaxation time-scale  $\tau_{\text{SBM}}$  of 2 hours and a refer-  
122 ence relative humidity  $RH_{\text{SBM}} = 0.7$ . The boundary layer scheme is the one used by O’Gorman  
123 and Schneider (2008). The moist GCM is run under perpetual equinox conditions, with no daily  
124 cycle of insolation, and is coupled to a slab ocean of depth 1m, with no representation of ocean

125 dynamics or of sea ice. A mixed-layer depth of 1m was used so that the model would spin up  
126 quickly; using a deeper mixed-layer damps the temperature variance, but otherwise our results  
127 are qualitatively insensitive to the choice of mixed-layer depth. Moreover, a mixed-layer depth of  
128 1m allows surface temperatures to respond to synoptic-scale forcing, as continental land surfaces  
129 do. A deeper mixed-layer depth, more representative of an oceanic mixed-layer, would decouple  
130 surface temperatures from synoptic temperature variability.

131 The same Gaussian orography is added to the model as in the dry GCM, except that it is centered  
132 further north at 60°N. The reason for moving the orography poleward is that the storm tracks, and  
133 the associated maxima in temperature variance, are further poleward in this set-up (see Figure 3  
134 below), so a more northward mountain produces clearer changes in variance. As discussed by  
135 Wills and Schneider (2018), this implementation of orography produces an “aqua-mountain”, and  
136 the surface fluxes over the orography are not necessarily realistic. However, any bias in the surface  
137 fluxes is of secondary importance for our investigation.

138 The moist GCM was integrated at T85 truncation with 30 unevenly-spaced vertical levels, start-  
139 ing from a state with uniform SSTs. The simulations lasted for 4500 days with data stored four  
140 times per day, and we have taken averages over the final 4000 days.

141 Our focus in this study is on winter temperature variability, as land surface processes, like soil-  
142 moisture feedbacks, are less important for variability in winter than in summer. As neither of the  
143 idealized GCMs includes a representation of land surface processes, they can be used to study the  
144 mechanisms of winter temperature variance without imposing seasonality and so, for convenience,  
145 we have used set-ups that produce equinoctial climates.

146 *d. Comprehensive climate model*

147 The comprehensive climate model is GFDL CM2.5-FLOR (Vecchi et al. 2014). FLOR stands for  
148 Forecast-oriented Low Ocean Resolution, and the model is based on the GFDL CM2.5 model. It  
149 is run with an atmospheric resolution of approximately 50km and an oceanic resolution of approx-  
150 imately  $1^\circ$ . By running with a relatively high resolution atmosphere, FLOR is able to accurately  
151 capture many subseasonal forms of variability, such as hurricanes and monsoon depressions, and  
152 can resolve sharp topographic features, such as the peaks of the Himalayas (compare panels a and  
153 b of Figure 2).

154 Three simulations were performed with FLOR: (1) a control simulation with present-day to-  
155 pography, (2) a simulation with the Rockies flattened to 300m (the “no-Rockies” simulation, i.e.,  
156 all surface heights greater than 300m are reduced to 300m) and (3) a simulation with the Asian  
157 orography (the Tibetan Plateau, the Himalayas and the Mongolian Plateau) flattened to 300m (the  
158 “no-Tibet” simulation). The regions of flattened topography can be seen in Figure 2 and we note  
159 that the gravity wave drag and boundary layer roughness were fixed to their control values where  
160 the topography was flattened (see also Baldwin et al. (2019b)).

161 All simulations were conducted with pre-industrial radiative forcings, matching the best guess  
162 for the year 1860, and with static vegetation. Daily-mean data were collected for 50 years, fol-  
163 lowing 100 years of spin-up from an initial state of rest, and SSTs were relaxed to a repeating  
164 climatology with a relaxation time-scale of five days. This set-up was originally designed to al-  
165 low tropical cyclones to interact with the ocean surface (Vecchi et al. 2014); for our purposes, the  
166 model is essentially an atmosphere-only climate model run over fixed SSTs. Our configuration  
167 attempts to isolate the direct effects of the orographic forcing on temperature variability, though  
168 not the indirect effects orography has on variability through its impact on SSTs.

169 *e. Filtering to synoptic time-scales*

170 The data were filtered to synoptic time-scales using a fourth-order Butterworth filter, with cut-  
171 off frequencies of  $1/3 \text{ days}^{-1}$  and  $1/15 \text{ days}^{-1}$ . The filter was implemented using the Python  
172 package `scipy.signal`, with the filter co-efficients obtained using `scipy.signal.butter` and  
173 the filtering done with `scipy.signal.lfilter`. We have verified that our results are robust to  
174 the choice of filtering time-scales, within reason. For all datasets, DJF variance and skewness were  
175 calculated individually for each year (e.g., December 1979 to February 1980) and then averaged  
176 over all years to find the climatological variance and skewness.

177 **3. Impact of Orography on Temperature Variance in Idealized Models**

178 *a. Background theory*

179 Assuming that synoptic potential temperature variations are primarily generated by horizontal  
180 advection, and that this advection is local in time and space, potential temperature variations can  
181 be Taylor expanded to give (Corrsin (1974); Schneider et al. (2015))

$$\theta' \approx -\frac{\partial \bar{\theta}}{\partial y} L'_y - \frac{\partial \bar{\theta}}{\partial x} L'_x + \frac{1}{2} \frac{\partial^2 \bar{\theta}}{\partial y^2} L'^2_y + \dots, \quad (2)$$

182 where  $\theta' = \theta - \bar{\theta}$  denotes synoptic variations of potential temperature at 850hPa about some lo-  
183 cal mean value  $\bar{\theta}$ ,  $L'_y$  is the Lagrangian displacement of air masses arriving at  $y$  from  $y_0$ , and  
184 similarly for  $L'_x$ . “Mean” denotes an average over a time-scale that is long compared to synoptic  
185 time-scales and we consider potential temperature rather than temperature because potential tem-  
186 perature is materially conserved during adiabatic air mass displacements. We work in Cartesian  
187 co-ordinates for simplicity, and define  $L'_y$  as positive for a northward displacement and  $L'_x$  as posi-  
188 tive for an eastward displacement. Provided the length scales of potential temperature variations,  
189  $\bar{L}_y = 2|\partial_y \bar{\theta} / \partial_{yy} \bar{\theta}|$  and  $\bar{L}_x = 2|\partial_x \bar{\theta} / \partial_{xx} \bar{\theta}|$ , are much larger than the mixing length-scales,  $L'_y$  and  $L'_x$ ,

190 the expansion can be well approximated by just retaining the first two terms, and so the synoptic  
 191 potential temperature variance can be approximated as

$$\overline{\theta'^2} \approx \overline{\left(\frac{\partial \bar{\theta}}{\partial y}\right)^2} L_y^2 + \overline{\left(\frac{\partial \bar{\theta}}{\partial x}\right)^2} L_x^2 + 2 \overline{\left(\frac{\partial \bar{\theta}}{\partial y}\right) L_y \left(\frac{\partial \bar{\theta}}{\partial x}\right) L_x}. \quad (3)$$

192 The meridional term, specifically the meridional temperature gradient, generally dominates over  
 193 the zonal term and the cross term (note the different colorbar scales in panels c and d of Figure 1),  
 194 but we have included the latter two here to emphasize that zonal temperature gradients also impact  
 195 regional potential temperature variability.

196 Orography affects temperature gradients by meridionally compressing downstream near-surface  
 197 isentropes and pulling apart upstream isentropes. However, this requires the flow to be deflected  
 198 around the orography, rather than flowing up and over it, so that the air deflected equatorward  
 199 partly adjusts to the warmer conditions and the air deflected poleward partly adjusts to the colder  
 200 conditions, before the downstream confluence of the flow. For small heights the air flows up and  
 201 over the orography, leaving the potential temperature gradients unaffected. Formally, consider the  
 202 linearized, time-mean thermodynamic equation for adiabatic flow on the lowest model level (in  $z$   
 203 coordinates):

$$\bar{u} \frac{\partial \theta'}{\partial x} + v' \frac{\partial \bar{\theta}}{\partial y} = -w \frac{\partial \bar{\theta}}{\partial z}. \quad (4)$$

204 The orographic forcing enters through the lower boundary condition, which can be approximated  
 205 in the linear regime as (Cook and Held 1992)

$$w \approx \bar{u} \frac{\partial h}{\partial x}, \quad (5)$$

206 where  $h$  is again the height of the orography, with maximum height  $H$ . Substituting then gives

$$\bar{u} \frac{\partial \theta'}{\partial x} + v' \frac{\partial \bar{\theta}}{\partial y} = -\bar{u} \frac{\partial h}{\partial x} \frac{\partial \bar{\theta}}{\partial z}. \quad (6)$$

207 In this linear regime, the air moves up and over the mountain, and the first term on the left side of  
 208 equation 6 balances the right side so that  $\frac{\theta'}{H} \sim \frac{\partial \bar{\theta}}{\partial z}$ . For larger  $H$  the flow becomes increasingly non-  
 209 linear, and the deflection around the orography is important. In this regime the forcings associated  
 210 with the meridional wind and with zonal wind anomalies can no longer be ignored in equation 5,  
 211 and the orographic forcing is balanced by the  $v' \frac{\partial \bar{\theta}}{\partial y}$  term<sup>1</sup>. In idealized experiments this transition  
 212 occurs for  $H$  between 1 and 2km for orography with approximately the same horizontal extent as  
 213 the Tibetan Plateau (Cook and Held (1992), Lutsko and Held (2016)). Valdes and Hoskins (1991)  
 214 demonstrated that Asian topography meets this criterion, but caution that it is less clear whether  
 215 the Rockies do, with the result depending on how the Rockies are defined. Furthermore, while  
 216 the near-surface flow appears to be deflected around the Rockies (see Figure 9 below), this flow  
 217 is strongly influenced by heating in the north Pacific storm track (see also Valdes and Hoskins  
 218 (1989)).

219 Another factor which enhances the downstream temperature gradients is the preferential deflec-  
 220 tion of the flow around the poleward side of the orography. If the flow follows isentropes, then  
 221 it will descend in height when it moves equatorward and ascend in height when it moves pole-  
 222 ward, following the mean isentropic slope (Valdes and Hoskins 1991). Thus the mountain appears  
 223 “taller” to the flow on its equatorward flank and “shorter” on its poleward flank, so that more of  
 224 the air flows around the poleward flank of the mountain. The downstream convergence is then  
 225 equatorward of the center of the orography, with anomalously cold air meeting the warm air that  
 226 flowed around the equatorward side of the orography.

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<sup>1</sup>Though note that the potential temperature perturbation is itself proportional to the deflection of the flow:  $\theta' \approx \eta' \frac{\partial \bar{\theta}}{\partial y}$ , where  $\eta$  is the typical meridional displacement of a fluid parcel, assumed to be equal to the meridional extent of the orography. Hence the condition for meridional deflection to dominate is  $|\eta/H| < \left| \frac{\partial \bar{\theta}}{\partial z} / \frac{\partial \bar{\theta}}{\partial y} \right|$ . Roughly, the meridional slope of the mountain must be greater than the characteristic slope of the isentropes (Valdes and Hoskins 1991).

227 *b. Idealized GCM results*

228 In both idealized GCMs, temperature variance is reduced upstream and enhanced downstream  
229 of orography (Figure 3a and b), as are meridional temperature gradients (panels c and d). However  
230 the inferred mixing lengths  $L' = \sqrt{\theta'^2 / \left(\frac{\partial \bar{\theta}}{\partial y}\right)^2}$  are reduced downstream of the orography (panels  
231 e and f), which is the result of two competing effects. First, by increasing downstream temperature  
232 gradients, orography increases downstream Eady growth rates, potentially leading to more ener-  
233 getic eddies and thus to larger mixing lengths (see Caballero and Hanley (2012) for discussion of  
234 the relationship between eddy kinetic energy and mixing lengths). But in addition to local baro-  
235 clinicity, eddies in strong jets are also energized by downstream development – by the recycling  
236 of energy from upstream eddies (Chang and Orlanski (1993); Chang et al. (2002)). Orography  
237 disrupts the latter by interfering with the zonal propagation of wave packets (Son et al. 2009), and  
238 for the set-ups used here this effect wins out, resulting in less energetic eddies and smaller effec-  
239 tive mixing lengths. This reduction in the mixing lengths has a substantial impact on the local  
240 response of the variance. For instance, panel a of Figure 4 shows the zonal anomalies in synoptic  
241 temperature variance for the simulation with the dry GCM and  $H = 4\text{km}$ , and it can be seen that  
242 the reduction in the mixing lengths creates a small region, near  $120^\circ\text{E}$ , in which the downstream  
243 variance is reduced, while the largest increase in variance is further downstream, at around  $170^\circ\text{E}$ ,  
244 where the eddies are more energetic.

245 On the poleward side of the mountain the pattern is reversed (Figure 4a), with enhanced temper-  
246 ature variance upstream and reduced variance downstream of the mountain. This is partly caused  
247 by the preferential deflection of the flow around the poleward flank of the mountain (arrows in Fig-  
248 ure 3a), which induces convergence on the northwest flank of the mountain, and thus a tightening  
249 of the isentropes, and divergence on the northeast flank of the mountain, causing the isentropes to



250 pull apart (see contours in Figure 3c). The jet is also relatively narrow in the dry GCM, compared  
251 to typical winter climates, so that there are strong polar easterlies at the latitudes of the poleward  
252 edge of the mountain. Hence the northeast flank is upstream of the mountain, and temperature  
253 variance should be reduced there.

254 Our focus is on the jet regions, however, where the enhanced meridional temperature gradients  
255 cause a local enhancement of temperature variance downstream of the orography in both models.  
256 Figure 5 shows that the maximum zonal anomaly in potential temperature variance increases in the  
257 simulations with the dry and moist GCMs as the height of the orography ( $H$ ) is increased (panel  
258 a)<sup>2</sup>, as does the maximum zonal anomaly of the squared meridional temperature gradient (panel  
259 b). Plotting these against each other demonstrates the strong linear relationship between the two  
260 quantities in the GCMs (panel c). The different slopes indicate that the mixing lengths differ in the  
261 two models, and the larger slope for the moist GCM implies that adding moist processes increases  
262 the effective mixing length (see below).

263 A possible complication is the shape of the orography: the Rockies form a meridionally-  
264 elongated ridge, whereas the Himalayas are more zonally-elongated. To investigate how the orog-  
265 raphy's shape influences temperature variability, two additional simulations were run with the dry  
266 GCM, one with a 4km meridional ridge resembling the Rockies ( $\alpha = 15^\circ$  and  $\beta = 5^\circ$ ) and one  
267 with a 4km zonal ridge ( $\alpha = 5^\circ$  and  $\beta = 15^\circ$ ).

268 Panels b and c of Figure 4 show the zonal anomalies in temperature patterns in these simulations  
269 (we note that the zonal-mean variance is lower in both of the ridge experiments than in the circular  
270 experiment because the ridges interfere less with the downstream development and hence the mix-  
271 ing lengths are larger than in the circular experiment). The zonal anomalies are broadly similar in

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<sup>2</sup>Note that the zonal-mean variance decreases with increasing  $H$  in both models because of the increasing disruption of downstream development by the orography (not shown).

272 all three experiments, with reductions in temperature variance upstream of the mountains and en-  
273 hancements downstream of the mountain and a reversed pattern at higher latitudes, however there  
274 are some noticeable differences. For instance, in the zonal ridge case the reduction is mostly on  
275 the southern flank of the mountain, rather than to the southwest. The meridional ridge produces a  
276 similar response to the circular experiment, but a key difference is that the variance is increased on  
277 the entire eastern flank of the meridional ridge. The Rockies show a similar local enhancement of  
278 variance on their eastern flank (Figure 1). In the meridional ridge simulation the largest increase  
279 in variance is also immediately downstream of the orography, on its southeastern flank, instead of  
280 being displaced further downstream, as for the circular case. Decomposing this response into a  
281 squared gradient and an inferred mixing length shows that in the meridional ridge case the tem-  
282 perature gradient is more strongly increased immediately downstream of the orography, relative to  
283 the reduction in the mixing length (not shown).

284 In the dry GCM, advection is the sole method of generating potential temperature variance,  
285 whereas in the moist GCM covariance of anomalous latent heating and potential temperature  
286 anomalies also contributes. To investigate the role of latent heat anomalies, Figure 6 shows the ad-  
287 vective terms in the temperature variance budget (see equation 3 of Wilson and Williams (2006))  
288 for the  $H = 4\text{km}$  simulation with the moist GCM, as well as the contribution of latent heat fluctu-  
289 ations to temperature variance ( $\overline{\theta'Q'_L}$ , Figure 6d). Latent heating enhances temperature variability  
290 downstream of the orography, increasing the inferred mixing lengths diagnosed in the moist GCM  
291 and partly explaining why the downstream maximum in temperature variability is closer to the  
292 mountain in this GCM than in the dry GCM. This enhancement is around 20% of the advective  
293 tendency, which is dominated by the  $\overline{\mathbf{u}'\theta' \cdot \nabla\bar{\theta}}$  term, and comes about because the latent heating  
294 fluctuations are related to the advection. For instance, anomalously warm air, originating close to

295 the surface in the tropics, will condense water as it moves poleward and rises, further enhancing  
296 the temperature anomalies.

297 In summary, the results of the GCM simulations agree with the theoretical expectations from  
298 the previous section, with reductions and enhancements of temperature variance caused mostly by  
299 changes in meridional temperature gradients due to the presence of orography. This is complicated,  
300 however, by reductions in the effective mixing lengths due to the interference of the orography with  
301 downstream development. The ridge experiments with the dry GCM also demonstrated important  
302 dependencies on the aspect ratio of the orography. In the case of a meridional ridge, resembling  
303 the Rockies, the variance is enhanced immediately downstream of the orography, whereas with  
304 a more “circular” orography the largest enhancement is further downstream. Finally, analyzing  
305 the temperature variance of the moist GCM demonstrates that the contribution of latent heating  
306 anomalies to temperature variance enhances the variance due to horizontal advection, as these  
307 latent heating anomalies are tied to the advection itself. So we can proceed by focusing on the  
308 advection, noting that latent heating enhances the effective mixing lengths.

#### 309 **4. Temperature Variability in Simulations with Flattened Orography**

310 Figure 7a shows that FLOR is able to reproduce the main features of MERRA’s pattern of DJF  
311 synoptic temperature variance<sup>3</sup>. In Figure 7b it can be seen that the effect of the Asian orography  
312 is to decrease the temperature variance over most of Eurasia, as well as over the North Pacific and  
313 North America, and to increase the variance over central Siberia (see also Figure 8). Notably, tem-  
314 perature variance is reduced over the heavily populated southeast Asian coast, including southern  
315 China, in the control simulation compared to the no-Tibet simulation. In part, this is because at

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<sup>3</sup>The temperature variance is somewhat higher in the FLOR simulations than in the reanalysis, which we attribute in part to the higher resolution of FLOR’s atmospheric model compared to the reanalysis: coarse-graining the data from the control simulation to a 1.25° grid reduces the synoptic temperature variance by about 30% on average (not shown). See also Supplementary Figure 3 of Baldwin et al. (2019a).

316 these latitudes the zonal-winds transition from westerly to easterly and this region is upstream of  
317 the orography (Figure 9). However, the primary cause of the reduced variance is the Asian orogra-  
318 phy's interference with downstream development, which weakens the storms over southeast Asia  
319 and, especially, in the Pacific storm track (Figure 10c). The Kuroshio Extension off the east coast  
320 of Japan is the genesis region for the Pacific storm track, and the Himalayas and Tibet weaken  
321 the eddies formed over the Kuroshio because of the reduced energy from upstream, despite the in-  
322 creased temperature gradient in the northwest Pacific. The reduced downstream development also  
323 impacts the strength of winter storms originating in the Pacific storm track and reaching North  
324 America.

325 The winter stationary wave pattern over Eurasia consists of a zonally oriented dipole, with  
326 anomalous warmth over Europe and anomalous cold over east Asia (Figure 7g). The presence  
327 of Tibet cools east Asia (compare Figure 7 panels g and h), implying that the stationary wave  
328 forced by the Asian orography constructively interferes with the stationary wave excited by the  
329 land-sea contrast on Eurasia's east coast (Kaspi and Schneider (2011); Park et al. (2013)). In the  
330 absence of the Asian orography the largest Eurasian temperature gradients are at relatively low  
331 latitudes, with the maximum gradient at about 30°N (Figure 7e), whereas the mid-latitude jet,  
332 where the mixing lengths are largest, is further north. This southward displacement of the maxi-  
333 mum temperature gradient when the orography is flattened contributes to the smaller temperature  
334 variance over Eurasia compared to North America in the no-Tibet simulation.

335 The Rockies act to increase the variance over most of North America, but also decrease the  
336 variance off the west coast of North America (panel c of Figure 7, Figure 8). Both the Rockies  
337 and the Asian orography increase the temperature variance over the polar regions, because their  
338 presence cools the high latitudes, increasing the zonal-mean equator-to-pole temperature gradient  
339 (Figure 10a). We have not fully diagnosed the reasons for this, but note that the mid-latitude jets

340 weaken in the presence of the mountain ranges, resulting in weaker poleward transient eddy heat  
341 fluxes (Figure 10 panels b and c).

342 Table 1 quantifies the changes in temperature variance over the two continents by comparing DJF  
343 synoptic temperature variance in the three FLOR simulations over a Eurasian box ( $40^{\circ}$ - $120^{\circ}$ E and  
344  $30^{\circ}$ - $75^{\circ}$ N) and over a North American box ( $240^{\circ}$ - $280^{\circ}$ E and  $30^{\circ}$ - $75^{\circ}$ N). The areas of the Asian  
345 mountains and the Rockies are masked whenever an average is taken over these boxes. Asian  
346 orography reduces the variance over the Eurasian box by  $1.4\text{K}^2$  and over the North American box  
347 by  $1.3\text{K}^2$ , with both of these changes statistically significant at the 95% level based on a two-sided  
348 Student's t-test. The Rockies enhance the variance over the North American box by  $1.3\text{K}^2$  and  
349 over Eurasia by  $0.2\text{K}^2$ , though only the change over North America is statistically significant in  
350 this case.

351 These calculations suggest that the enhancement of North American temperature variability by  
352 the Rockies is roughly canceled by the damping of variability due to Asian orography. The in-  
353 creases and decreases in variance are sensitive to the definitions of the boxes, however, and this  
354 cancellation also assumes the effects of flattening the mountain ranges individually can be linearly  
355 added together. Regardless, the majority of the orography's net effect comes from the reduction  
356 of Eurasian temperature variability by the Asian mountains and, in FLOR, this explains about a  
357 quarter of the difference in variance over the two continents ( $1.4\text{K}^2 / 5.5\text{K}^2 \approx 25\%$ ).

358 Our framework for explaining differences in temperature variance is based on differences in  
359 mean temperature gradients, which are in turn controlled by the Northern Hemisphere stationary  
360 wave pattern. So the remaining difference in temperature variance between the two continents can  
361 largely be attributed to stationary waves forced by diabatic heating, which, together with orography  
362 are responsible for the bulk of the Northern Hemisphere stationary wave pattern (Held et al. 2002).

363 Even in the no-Rockies simulation there are substantial meridional temperature gradients over  
364 North America (Figure 7f), and the stationary wave pattern over North America is similar in all  
365 three simulations, consisting of a dipole with anomalously warm temperatures off the west coast of  
366 North America and anomalously cold temperatures centered over northeast Canada (Figure 7 pan-  
367 els g), h) and i)). The dipole is weaker in the no-Rockies simulation, indicating that the stationary  
368 wave forced by the Rockies constructively interferes with the dipole. In this case, the pattern over  
369 North America is a combination of the stationary wave forced by the land-sea contrast between  
370 the east coast of North America and the western Atlantic (Kaspi and Schneider 2011), which cools  
371 eastern North America, and stationary waves forced by diabatic heating in the Pacific warm pool  
372 region and by thermal forcing in the extratropical Pacific (Hoskins and Karoly (1981); Valdes and  
373 Hoskins (1991); Held et al. (2002)). The latter includes the forcing due to the warm waters of the  
374 Kuroshio as well as the eddy sensible heat flux convergence in the Pacific storm track, making it  
375 difficult to separate out the relative contributions of the different thermal forcings.

## 376 **5. Temperature Skewness**

377 Through its effects on temperature gradients, orography also impacts the skewness of synoptic  
378 temperatures. Garfinkel and Harnik (2017) showed that, in mid-latitudes, synoptic temperature  
379 extremes occur when air is advected over regions with large mean meridional temperature gradi-  
380 ents, so that temperatures poleward of these regions tend to be positively skewed and temperatures  
381 equatorward of these regions tend to be negatively skewed. By strengthening downstream temper-  
382 ature gradients, orography increases the positive skewness to the north of these gradients and the  
383 negative skewness to the south. This is illustrated in Figure 11, which shows maps of skewness  
384 in simulations with the two idealized GCMs, as well as the meridional temperature gradients. In

385 both cases, downstream temperatures are skewed more positively north of the enhanced tempera-  
386 ture gradients and more negatively to the south of the gradients.

387 In the reanalysis data, the strongest DJF meridional temperature gradients are found in the storm  
388 track regions of the west Pacific and the west Atlantic (Figure 1c). Panel a of Figure 12 shows that  
389 synoptic temperatures are positively skewed in the northwest Pacific and the northwest Atlantic,  
390 and negatively skewed to the south of these regions. The same patterns are seen in the control  
391 simulation with FLOR (Figure 12b, note that as with the variance, we attribute the larger values  
392 of skewness in part to FLOR's higher resolution). The temperature gradient in the west Pacific is  
393 reduced in the no-Tibet simulation, and comparing panels b and c of Figure 12 confirms that the  
394 skewness in the northwest Pacific is also reduced in this simulation. Averaging over the region  
395 35°N-50°N and 140°E to 180°E (green box in Figure 12b) gives a reduction in skewness of 31%  
396 ( $= (0.234 - 0.162) / 0.234$ , difference significant at the 90% level) in the northwest Pacific.

397 Flattening the Rockies does not appear to affect temperature gradients in the west Atlantic (Fig-  
398 ure 7f), and the skewness in the northwest Atlantic is comparable in the control and the no-Rockies  
399 simulations. Over land, DJF synoptic temperatures are negatively skewed at almost all latitudes,  
400 and other factors, such as land-surface feedbacks, are likely important for generating extreme  
401 events.

## 402 **6. Conclusion**

403 In this study we have investigated the contribution of large-scale orography to the increased win-  
404 tertime near-surface daily and synoptic temperature variability over North America compared to  
405 Eurasia. Our analysis combines theoretical arguments, simulations with two idealized GCMs and  
406 simulations with a comprehensive climate model – GFDL CM2.5-FLOR – in which the Rockies  
407 and the Asian orography are separately flattened. These allow us to quantify the impacts these

408 mountain ranges have on temperature variability over North America and Eurasia, and suggest  
409 that large-scale Northern Hemisphere orography is responsible for roughly 25% of the difference  
410 in variability.

411 Large-scale orography enhances downstream temperature variability by meridionally compress-  
412 ing downstream isentropes and reduces upstream temperature variability because upstream isen-  
413 tropes are pulled apart. At the same time, the preferential deflection of the flow towards the pole-  
414 ward flank of the orography, together with the presence of high latitude easterlies, can cause this  
415 pattern to be reversed at high latitudes, with enhanced variance on the northwest flank and reduced  
416 variance on the northeast flank of the orography (in the Northern Hemisphere). We have also  
417 shown that the orography's aspect ratio can cause substantial differences in the pattern of variabil-  
418 ity; for instance, a meridional ridge, resembling the Rockies, induces a stronger local enhancement  
419 of temperature variance on its downstream flank, whereas for circular orography the enhanced  
420 variance is further downstream. Finally, latent heat anomalies reinforce temperature anomalies  
421 created by advection, as anomalously warm air originating from low latitudes condenses water as  
422 it moves poleward and rises.

423 Most of North America is downstream of the Rockies, so wintertime temperature variability  
424 is enhanced there, while the Asian orography is on the eastern edge of Eurasia, so temperature  
425 variability is damped over most of Eurasia. An important exception is the southeast Asian littoral,  
426 which is east of the orography but exhibits reduced temperature variability due to the Asian moun-  
427 tains. This is partly because these regions are at latitudes of mean easterlies, or in the transition  
428 from mean westerlies to mean easterlies, and hence are upstream of the Himalayas. Another fac-  
429 tor is interference by the Asian orography with the energization of eddies over the Asian continent  
430 and the Pacific storm track by downstream development. This results in weaker winter storms  
431 and reduced variability over the east Asian coast, the Pacific and North America. The reduction



432 in variability over North America due to the presence of the Asian orography is approximately as  
433 large as the increase due to the presence of the Rockies.

434 Orography also enhances downstream skewness, as regions to the north of the enhanced tem-  
435 perature gradient have more positively skewed temperatures and regions to the south have more  
436 negatively skewed temperatures. In the FLOR simulations, the Himalayas and the Tibetan Plateau  
437 are found to increase temperature skewness in the northwest Pacific by about 30%.

438 The remaining difference in synoptic temperature variability over North America compared to  
439 Eurasia is primarily due to a combination of diabatic heating in the Pacific warm pool region, air-  
440 sea fluxes over the warm Kuroshio current and eddy sensible heat flux convergence in the Pacific  
441 storm track (Valdes and Hoskins (1989); Held et al. (2002)). The smaller width of the North  
442 American continent and its northwest-southeast sloping western coastline may also be important  
443 – Brayshaw et al. (2009) explored how this influences the North Atlantic storm track. Separating  
444 out these different factors, and the non-linear interactions between them, is an important next step.

445 The dominant control of horizontal advection on winter synoptic temperature variability is a  
446 powerful tool for understanding the regional pattern of temperature variability, in today's climate  
447 and how it may change in the future. This simplifies the problem to understanding the boreal winter  
448 stationary wave pattern, for which there is a large body of literature that can be drawn on (e.g.,  
449 Hoskins and Karoly (1981); Held (1983); Held et al. (2002)), though differences in mixing lengths,  
450 for instance due to orographic interference with downstream development, are an important caveat.  
451 Similarly, past and future changes in temperature variability can potentially be tied to changes in  
452 the stationary wave pattern (see e.g., Löfverström et al. (2014) and Simpson et al. (2016) for  
453 investigations of past and future changes in Northern Hemisphere stationary waves). More work  
454 is needed to better understand the impact of orography on mixing lengths, as well as to account  
455 for land surface processes such as soil-moisture, which affect temperature variability, particularly

456 during summer. These factors are also important for temperature extremes, particularly over land,  
457 where winter temperatures at almost all latitudes are negatively skewed. Nevertheless, the basic  
458 dynamics we describe here are robustly seen in idealized GCMs and in comprehensive climate  
459 models, and provide an important first step in explaining why North America experiences more  
460 wintertime temperature variability than Eurasia.

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## 467 **References**

- 468 Baldwin, J. W., J. B. Dessy, G. A. Vecchi, and M. Oppenheimer, 2019a: Temporally compound  
469 heat wave events and global warming: An emerging hazard. *Earth’s Future*, **7**.
- 470 Baldwin, J. W., G. A. Vecchi, and S. Bordoni, 2019b: The ocean-mediated influence of asian  
471 orography on tropical precipitation and cyclones. *Climate Dynamics*, 1–20.
- 472 Barnes, E. A., and L. M. Polvani, 2015: Cmp5 projections of arctic amplification, of the north  
473 american/north atlantic circulation, and of their relationship. *Journal of Climate*, **28 (3)**, 5254–  
474 5271.
- 475 Bathiany, S., V. Dakos, M. Scheffer, and T. M. Lenton, 2018: Climate models predict increasing  
476 temperature variability in poor countries. *Science Advances*, **4 (5)**.

- 477 Bowers, E. K., J. L. Grindstaff, S. S. Soukup, N. E. Drilling, K. P. Eckerle, S. K. Sakaluk, and C. F.  
478 Thompson, 2016: Spring temperatures influence selection on breeding date and the potential for  
479 phenological mismatch in a migratory bird. *Ecology*, **97**, 2880–2891.
- 480 Brayshaw, D. J., B. Hoskins, and M. Blackburn, 2009: The basic ingredients of the north at-  
481 lantic storm track. part i: Landsea contrast and orography. *Journal of the Atmospheric Sciences*,  
482 **66 (23)**, 2539–2558.
- 483 Caballero, R., and J. Hanley, 2012: Midlatitude eddies, storm-track diffusivity, and poleward  
484 moisture transport in warm climates. *Journal of the Atmospheric Sciences*, **69 (630)**, 3237–  
485 3250.
- 486 Chang, E. K. M., S. Lee, and K. L. Swanson, 2002: Storm track dynamics. *Journal of Climate*,  
487 **15 (16)**, 2163–2183.
- 488 Chang, E. K. M., and I. Orlanski, 1993: On the dynamics of a storm track. *Journal of the Atmo-  
489 spheric Sciences*, **50 (56)**, 999–1015.
- 490 Cohen, J., and Coauthors, 2014: Recent arctic amplification and extreme mid-latitude weather.  
491 *Nature Geoscience*, **7**, 627–637.
- 492 Cook, K. H., and I. M. Held, 1992: The stationary response to large-scale orography in a general  
493 circulation model and a linear model. *Journal of the Atmospheric Sciences*, **49 (6)**, 525–539.
- 494 Corrsin, S., 1974: Limitations of gradient transport models in random walks and in turbulence.  
495 *Advances in Geophysics*, **18A**, 25–60.
- 496 Frierson, D. M. W., 2007: The dynamics of idealized convection schemes and their effect on the  
497 zonally averaged tropical circulation. *Journal of the Atmospheric Sciences*, **64 (23)**, 1959–1976.

- 498 Frierson, D. M. W., I. M. Held, and P. Zurita-Gotor, 2006: A gray-radiation aquaplanet moist gcm.  
499 part i: Static stability and eddy scales. *Journal of the Atmospheric Sciences*, **63 (23)**, 2548–2566.
- 500 Garfinkel, C. I., and N. Harnik, 2017: The non-gaussianity and spatial asymmetry of temperature  
501 extremes relative to the storm track: The role of horizontal advection. *Journal of Climate*, **30**,  
502 445–464.
- 503 Held, I. M., 1983: Stationary and quasi-stationary eddies in the extratropical troposphere: theory.  
504 *Large-scale Dynamical Processes in the Atmosphere*, B. J. Hoskins, and R. P. Pearce, Eds.,  
505 Academic Press, New York, NY.
- 506 Held, I. M., and M. J. Suarez, 1994: A proposal for the intercomparison of the dynamical cores  
507 of atmospheric general circulation models. *Bulletin of the American Meteorological Society*,  
508 **75 (10)**, 1825–1830.
- 509 Held, I. M., M. Ting, and H. Wang, 2002: Northern winter stationary waves: Theory and modeling.  
510 *Journal of Climate*, **15 (16)**, 2125–2144.
- 511 Holmes, C. R., T. Woollings, E. Hawkins, and H. de Vries, 2016: Robust future changes in tem-  
512 perature variability under greenhouse gas forcing and the relationship with thermal advection.  
513 *Journal of Climate*, **29 (29)**, 2221–2235.
- 514 Hoskins, B., and T. Woollings, 2015: Persistent extratropical regimes and climate extremes. *Cur-  
515 rent Climate Change Reports*, **1 (3)**, 115–124.
- 516 Hoskins, B. J., and D. J. Karoly, 1981: The steady linear response of a spherical atmosphere to  
517 thermal and orographic forcing. *Journal of the Atmospheric Sciences*, **38 (6)**, 1179–1196.
- 518 Huybers, P., K. A. McKinnon, A. Rhines, and M. Tingley, 2014: U.s. daily temperatures: The  
519 meaning of extremes in the context of nonnormality. *Journal of Climate*, **27**, 7368–7384.

520 Jackson, S. T., J. L. Betancourt, R. K. Booth, and S. T. Gray, 2009: Ecology and the ratchet  
521 of events: Climate variability, niche dimensions, and species distributions. *Proceedings of the*  
522 *National Academy of Sciences*, **106**, 19 685–19 692.

523 Jahn, M., 2015: Economics of extreme weather events: Terminology and regional impact models.  
524 *Weather and Climate Extremes*, **10**, 29 – 39.

525 Kaspi, Y., and T. Schneider, 2011: Winter cold of eastern continental boundaries induced by warm  
526 ocean waters. *Nature*, **472**, 621–624.

527 Lau, N.-C., and M. J. Nath, 2012: A model study of heat waves over north america: Meteorological  
528 aspects and projections for the twenty-first century. *Journal of Climate*, **25**, 4761–4784.

529 Lau, N.-C., and M. J. Nath, 2014: Model simulation and projection of european heat waves in  
530 present-day and future climates. *Journal of Climate*, **27**, 3713–3730.

531 Lazo, J. K., M. Lawson, P. H. Larsen, and D. M. Walkman, 2011: U.s. economic sensitivity to  
532 weather variability. *Bulletin of the American Meteorological Society*, **92 (29)**, 709–720.

533 Linz, M., G. Chen, and Z. Hu, 2018: Large-scale atmospheric control on non-gaussian tails of  
534 midlatitude temperature distributions. *Geophysical Research Letters*, **45 (15)**, 9141–9149.

535 Löffverström, M., R. Caballero, J. Nilsson, and J. Kleman, 2014: Evolution of the large-scale  
536 atmospheric circulation in response to changing ice sheets over the last glacial cycle. *Climate of*  
537 *the Past*, **10 (4)**, 1453–1471.

538 Lutsko, N. J., and I. M. Held, 2016: The response of an idealized atmosphere to orographic  
539 forcing: Zonal versus meridional propagation. *Journal of the Atmospheric Sciences*, **73 (9)**,  
540 3701–3718.

541 McKinnon, K. A., A. Rhines, M. Tingley, and P. Huybers, 2016: The changing shape of northern  
542 hemisphere summer temperature distributions. *Journal of Geophysical Research: Atmospheres*,  
543 **121**, 8849–8868.

544 O’Gorman, P. A., and T. Schneider, 2008: The hydrological cycle over a wide range of climates  
545 simulated with an idealized gcm. *Journal of Climate*, **21 (15)**, 3815–3832.

546 Park, H.-S., S.-P. Xie, and S.-W. Son, 2013: Poleward stationary eddy heat transport by the tibetan  
547 plateau and equatorward shift of westerlies during northern winter. *Journal of the Atmospheric*  
548 *Sciences*, **70 (23)**, 3288–3301.

549 Rahmstorf, S., and D. Coumou, 2011: Increase of extreme events in a warming world. *Proceedings*  
550 *of the National Academy of Sciences*, **108 (44)**, 17 905–17 909.

551 Rienecker, M. M., and Coauthors, 2011: Merra: Nasa’s modern-era retrospective analysis for  
552 research and applications. *Journal of Climate*, **24**, 3624–3648.

553 Schneider, T., T. Bischoff, and H. Plotka, 2015: Physics of changes in synoptic midlatitude tem-  
554 perature variability. *Journal of Climate*, **28 (6)**, 2312–2331.

555 Screen, J. A., 2014: Arctic amplification decreases temperature variance in northern mid- to high-  
556 latitudes. *Nature Climate Change*, **4**, 577–582.

557 Shi, L., I. Kloog, A. Zanobetti, P. Liu, and J. D. Schwartz, 2015: Impacts of temperature and its  
558 variability on mortality in new england. *Nature Climate Change*, **5**, 988–991.

559 Simpson, I. R., R. Seager, M. Ting, and T. A. Shaw, 2016: Causes of change in northern hemi-  
560 sphere winter meridional winds and regional hydroclimate. *Nature Climate Change*, **6**, 65–70.

- 561 Son, S.-W., M. Ting, and L. M. Polvani, 2009: The effect of topography on storm-track intensity  
562 in a relatively simple general circulation model. *Journal of the Atmospheric Sciences*, **66 (23)**,  
563 393–311.
- 564 Valdes, P. J., and B. J. Hoskins, 1989: Linear stationary wave simulations of the time-mean clima-  
565 tological flow. *Journal of the Atmospheric Sciences*, **46 (18)**, 2509–2527.
- 566 Valdes, P. J., and B. J. Hoskins, 1991: Nonlinear orographically forced planetary waves. *Journal*  
567 *of the Atmospheric Sciences*, **48 (18)**, 2089–2106.
- 568 Vecchi, G. A., and Coauthors, 2014: On the seasonal forecasting of regional tropical cyclone  
569 activity. *Journal of Climate*, **27 (1)**, 7994–8016.
- 570 Wheeler, T., and J. von Braun, 2013: Climate change impacts on global food security. *Science*,  
571 **341 (6145)**, 508–513.
- 572 Wills, R. C., and T. Schneider, 2018: Mechanisms setting the strength of orographic rossby waves  
573 across a wide range of climates in a moist idealized gcm. *Journal of Climate*, **31 (18)**, 7679–  
574 7700.
- 575 Wilson, C., and R. G. Williams, 2006: When are eddy tracer fluxes directed downgradient? *Jour-*  
576 *nal of Physical Oceanography*, **36 (23)**, 189–202.

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580 and 30°N-75°N) in the FLOR simulations and observed variances from 1979-  
581 2012. All units are K<sup>2</sup> and the plus/minus values show the standard deviations  
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583 TABLE 1. Variance of December-January-February (DJF) 850hPa synoptic temperature over Eurasia (40°E-  
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 585 served variances from 1979-2012. All units are  $K^2$  and the plus/minus values show the standard deviations of  
 586 the interannual variability.

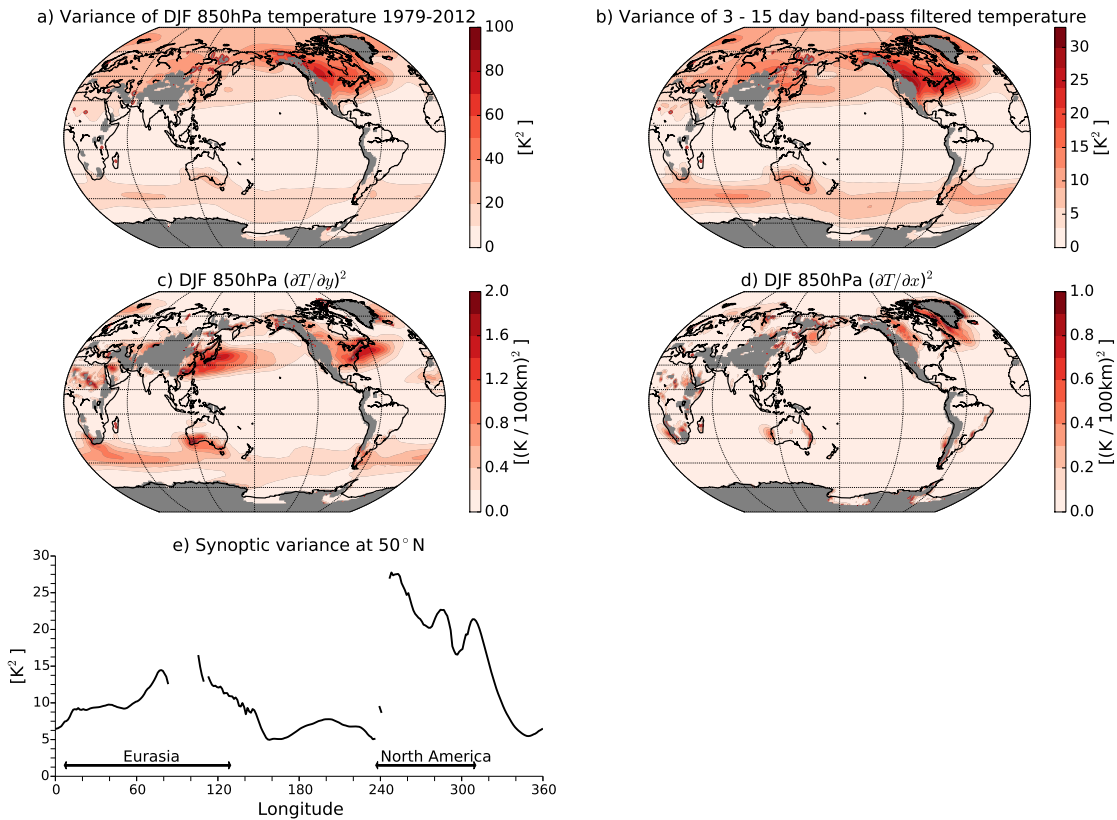
	MERRA reanalysis 1979-2012	Control	no-Tibet	no-Rockies
North America	$15.7 \pm 2.6$	$18.1 \pm 2.4$	$19.4 \pm 2.7$	$16.8 \pm 2.9$
Eurasia	$8.6 \pm 2.0$	$12.6 \pm 2.3$	$14.0 \pm 2.8$	$12.4 \pm 2.2$
Difference	7.1	5.5	5.4	4.4

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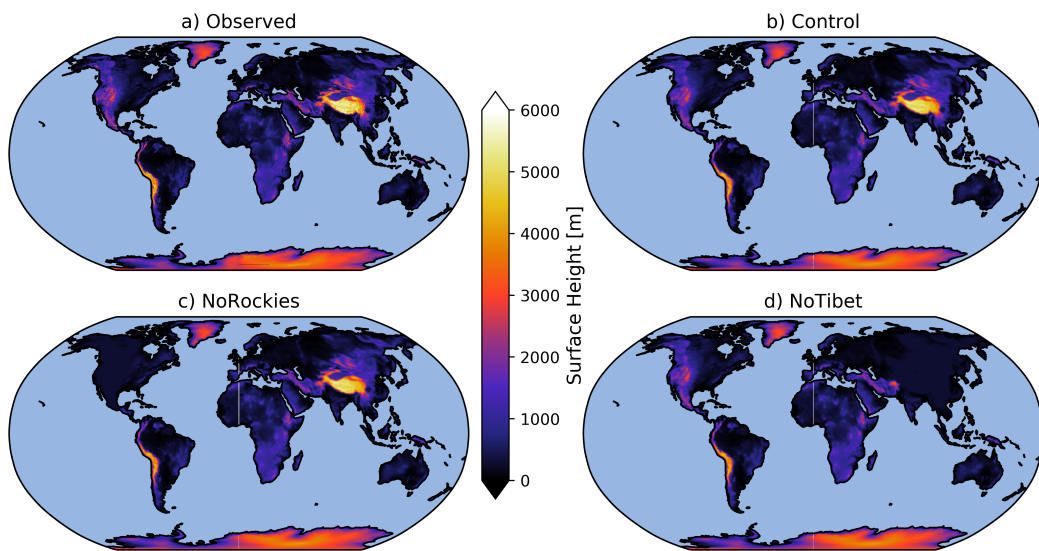
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- Fig. 6.** Panels a), b) and c): advective terms in the 850hPa potential temperature variance budget from a simulation with the moist GCM and a mountain height of 4km. Locations where topography intrudes through 850hPa are masked in gray. Panel d): the contribution of latent heating fluctuations to 850hPa potential temperature variance in the same simulation. . . . . 37
- Fig. 7.** a) Synoptic-scale variance of DJF 850hPa potential temperature in the control simulation with the comprehensive climate model, FLOR. b) Difference in synoptic-scale variance between the control simulation and the no-Tibet simulation. c) Difference in synoptic-scale

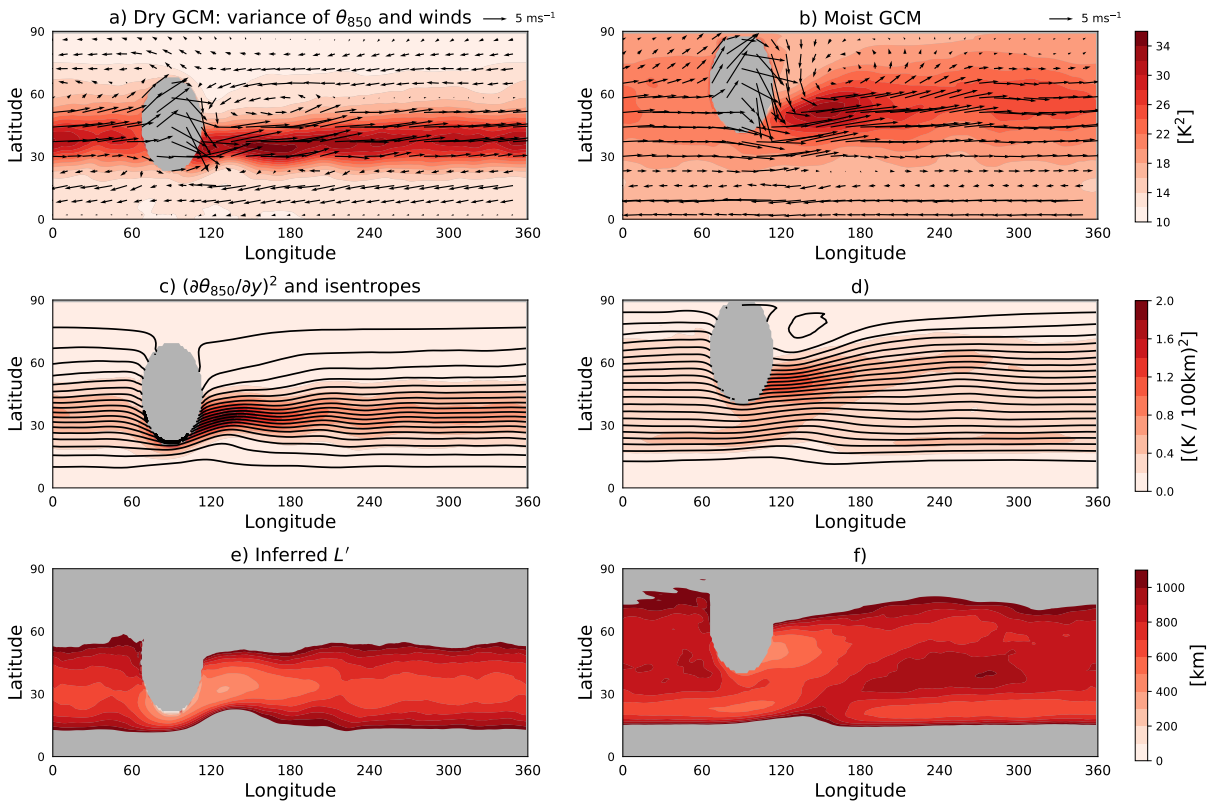
632	variance between the control simulation and the no-Rockies simulation. d) DJF squared	
633	meridional potential temperature gradients in the control simulation. e) DJF squared meridional	
634	potential temperature gradients in the no-Tibet simulation. f) DJF squared meridional	
635	potential temperature gradients in the no-Rockies simulation. g), h), i) DJF zonal anomalies	
636	in 850hPa potential temperature in the same simulations. Locations where topography	
637	intrudes through 850hPa are masked in gray. . . . .	38
638	<b>Fig. 8.</b> a) Profiles taken at 35°N of synoptic-scale variance of DJF 850hPa potential temperature in	
639	the three simulations with FLOR. b) Profiles taken at 50°N. Gaps in the profiles show where	
640	topography intrudes into the 850hPa level. . . . .	39
641	<b>Fig. 9.</b> a) DJF 850hPa temperature (contours) and total wind vectors in the vicinity of the Tibetan	
642	Plateau, averaged over the period 1979 to 2012. Data are taken from the MERRA reanalysis	
643	dataset. Locations where topography intrudes through 850hPa are masked in gray. b) Same	
644	for the region near the Rocky Mountains. . . . .	40
645	<b>Fig. 10.</b> a) Difference in zonal-mean $\theta_{850}$ between the control simulation with FLOR and the no-	
646	Tibet simulation (solid line) and difference between the control simulation and the no-	
647	Rockies simulation (dashed line). b) Differences in transient eddy potential temperature	
648	flux in the same simulations. c) Difference in DJF synoptic 850hPa eddy kinetic energy	
649	( $\overline{v'^2}$ ) between the control simulation and the no-Tibet simulation. . . . .	41
650	<b>Fig. 11.</b> a) Skewness of 850hPa synoptic temperatures (colored contours) and 850hPa meridional	
651	temperature gradients (black contours, contour interval = $0.2\text{K}(100\text{km})^{-1}$ ) in the dry GCM	
652	simulation with $H = 4\text{km}$ . b) Same for the simulation with the moist GCM. The meridional	
653	gradient contour interval is $0.2\text{K}/100\text{km}$ in both panels. . . . .	42
654	<b>Fig. 12.</b> a) Skewness of DJF 850hPa synoptic temperatures for the period 1979-2012 in the MERRA	
655	data. b) Skewness of DJF 850hPa synoptic temperatures in the control simulation with	
656	FLOR. c) Skewness of DJF 850hPa synoptic temperatures in the no-Tibet simulation with	
657	FLOR. d) Skewness of DJF 850hPa synoptic temperatures in the no-Rockies simulation with	
658	FLOR. . . . .	43



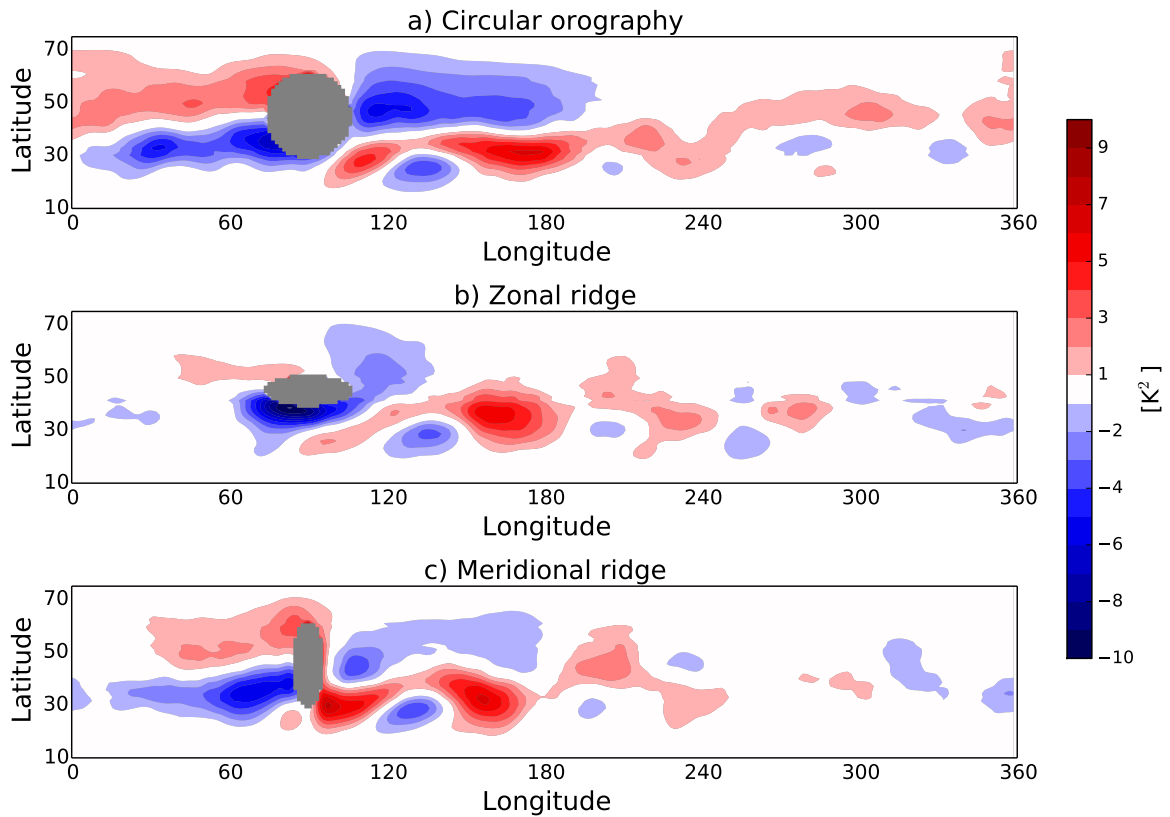
659 FIG. 1. a) Variance of December-January-February (DJF) 850hPa daily temperature for the period 1979 to  
 660 2012, calculated using data taken from the MERRA reanalysis dataset. Locations where topography intrudes  
 661 through 850hPa are masked in gray. b) Same as panel a, but the data are filtered using a fourth-order Butterworth  
 662 filter to only retain power at synoptic time-scales, here defined as 3 to 15 days. c) Climatological DJF squared  
 663 meridional temperature gradients for the same data. d) Climatological DJF squared zonal temperature gradients  
 664 for the same data. e) Profile of synoptic-scale 850hPa temperature variance at 50°N. Gaps in the profiles show  
 665 where topography intrudes into the 850hPa level.



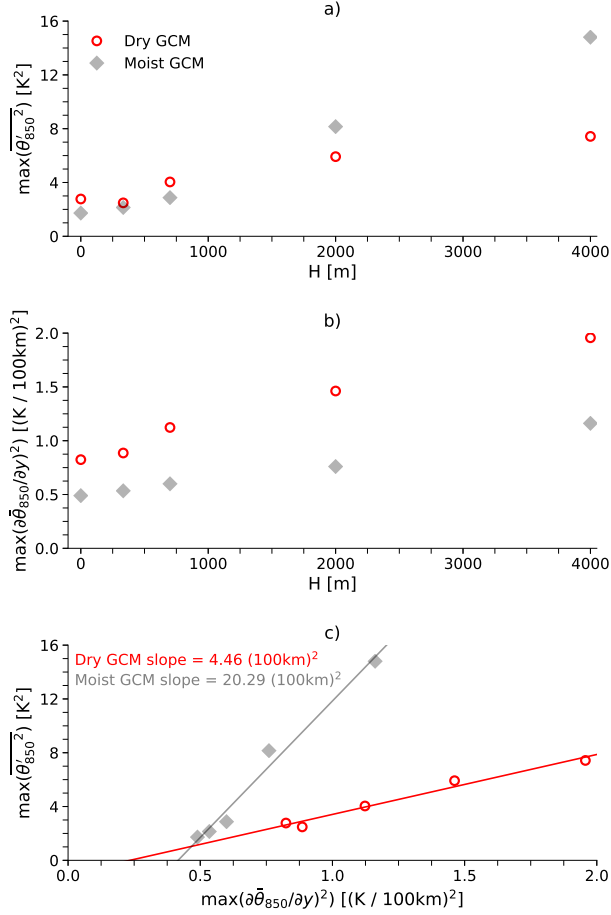
666 FIG. 2. a) Observed topography of Earth, taken from the ETOPO5 dataset, with 5 minute resolution. b)  
667 Topography in the control simulation of FLOR. c) Topography in the no-Rockies simulation. d) Topography in  
668 the no-Tibet simulation.



669 FIG. 3. a) Variance of synoptic (3 to 15 days) 850hPa potential temperature (contours) and total wind vectors  
 670 (arrows) in a simulation with the dry GCM and maximum orographic height of 4km. b) Same as panel a) but  
 671 for the simulation with the moist GCM with 4km orography. c) Squared meridional gradient of time-averaged  
 672 potential temperature (colored contours) and isentropes (black contours, with contour interval 2K) at 850hPa for  
 673 the same simulation as in panel a). d) Same as panel c) but for the simulation with the moist GCM with 4km  
 674 orography. e) Inferred mixing length  $L'$  for the simulation with the dry GCM and maximum orographic height  
 675 of 4km. f) Same as panel d) but for the simulation with the moist GCM with 4km orography. In all panels gray  
 676 indicates locations with surface pressure less than 850hPa or, in the bottom panels, where values are outside  
 677 the colorbar range. In a) and b) the winds are taken from the 0.85- $\sigma$  level so that the flow over and around the  
 678 orography is visible.

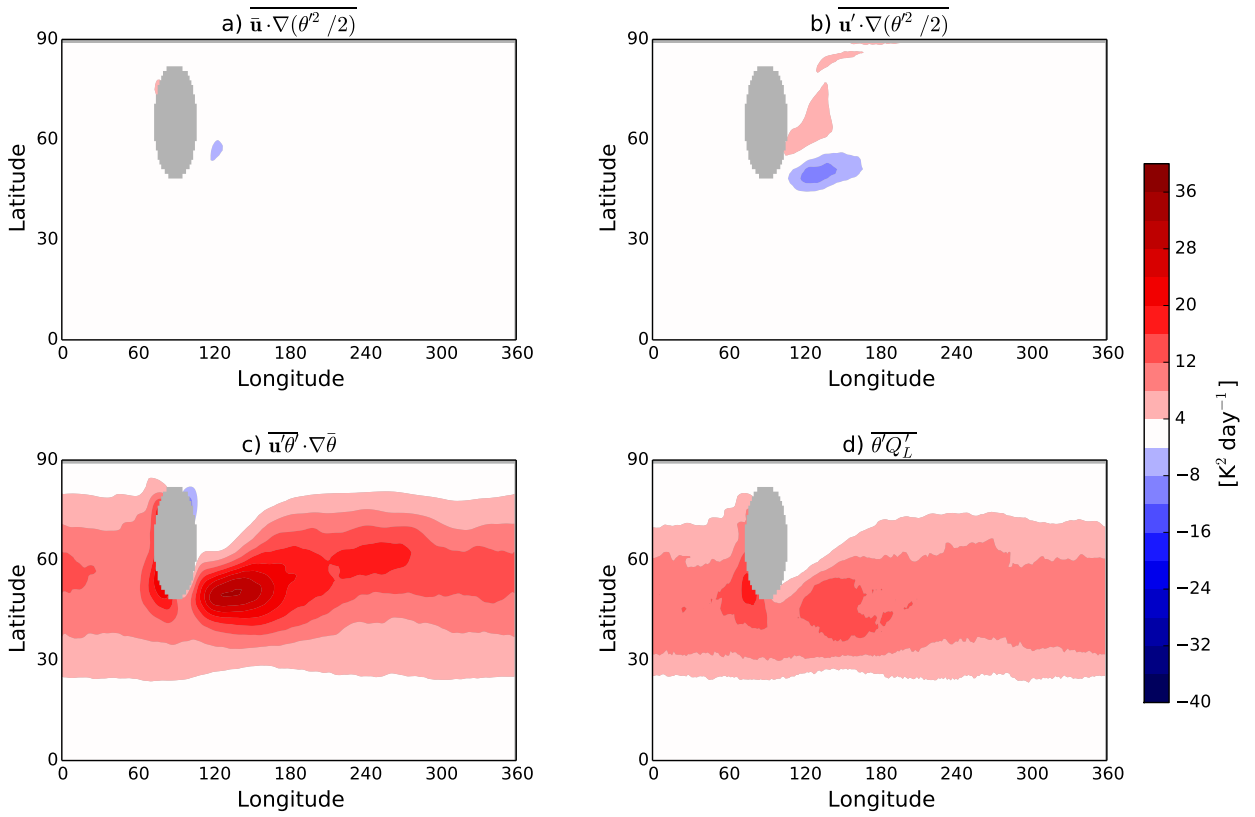


679 FIG. 4. a), b), c) Zonal anomalies in the variance of synoptic (3-15 day) 850hPa potential temperature in the  
 680 dry GCM simulations with  $H = 4\text{km}$  and the circular Gaussian orography (a,  $\alpha = \beta = 15^\circ$ ), the zonal ridge (b,  
 681  $\alpha = 5^\circ$  and  $\beta = 15^\circ$ ) and the meridional ridge (c,  $\alpha = 15^\circ$  and  $\beta = 5^\circ$ ).

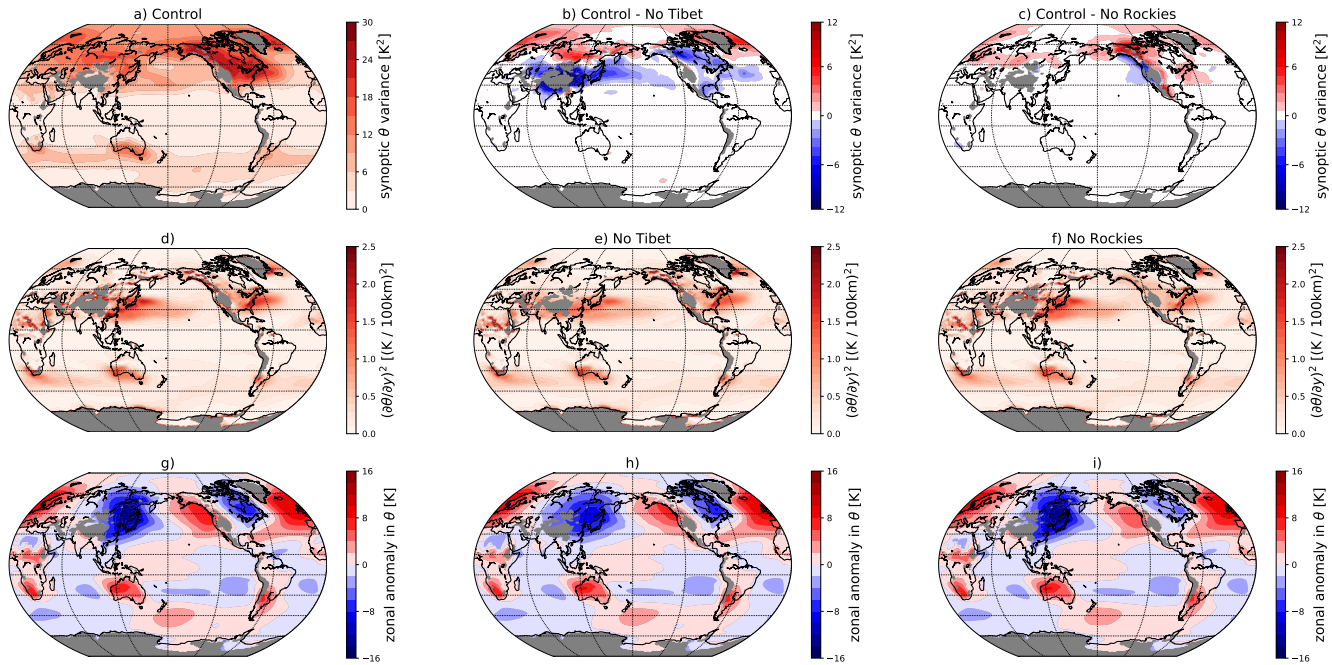


682 FIG. 5. a) Maximum anomalous 850hPa potential temperature variance ( $\max(\overline{\theta_{850}^2})$ ) as a function of moun-  
683 tain height,  $H$ , in the simulations with the dry GCM (red circles) and with the moist GCM (gray diamonds).  
684  $\max(\overline{\theta_{850}^2})$  is calculated as the maximum zonal anomaly in 850hPa potential temperature variance in the 120°  
685 downstream of the peak of the orography. b) Maximum zonal anomaly of the squared meridional potential tem-  
686 perature gradient at 850hPa ( $\max(\overline{(\partial \bar{\theta}_{850} / \partial y)^2})$ ) as a function of  $H$  in the simulations with the idealized GCMs.  
687  $\max(\overline{(\partial \bar{\theta}_{850} / \partial y)^2})$  is calculated as the maximum zonal anomaly in the 850hPa meridional potential temperature  
688 gradient in the 120° downstream of the peak of the orography. c)  $\max(\overline{\theta_{850}^2})$  versus  $\max(\overline{(\partial \bar{\theta}_{850} / \partial y)^2})$  in the  
689 simulations with the idealized GCMs. The lines show linear least-squares fits to the two sets of simulations.

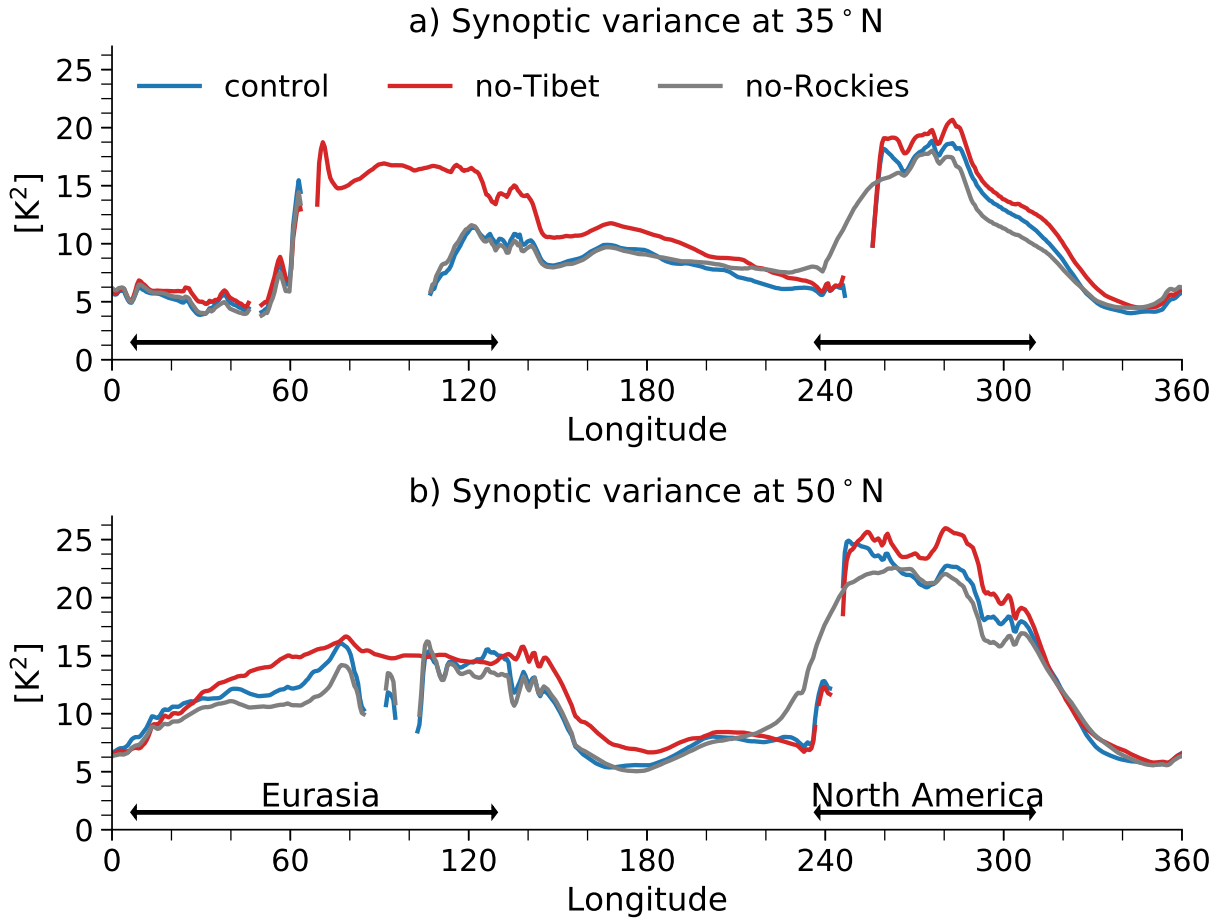




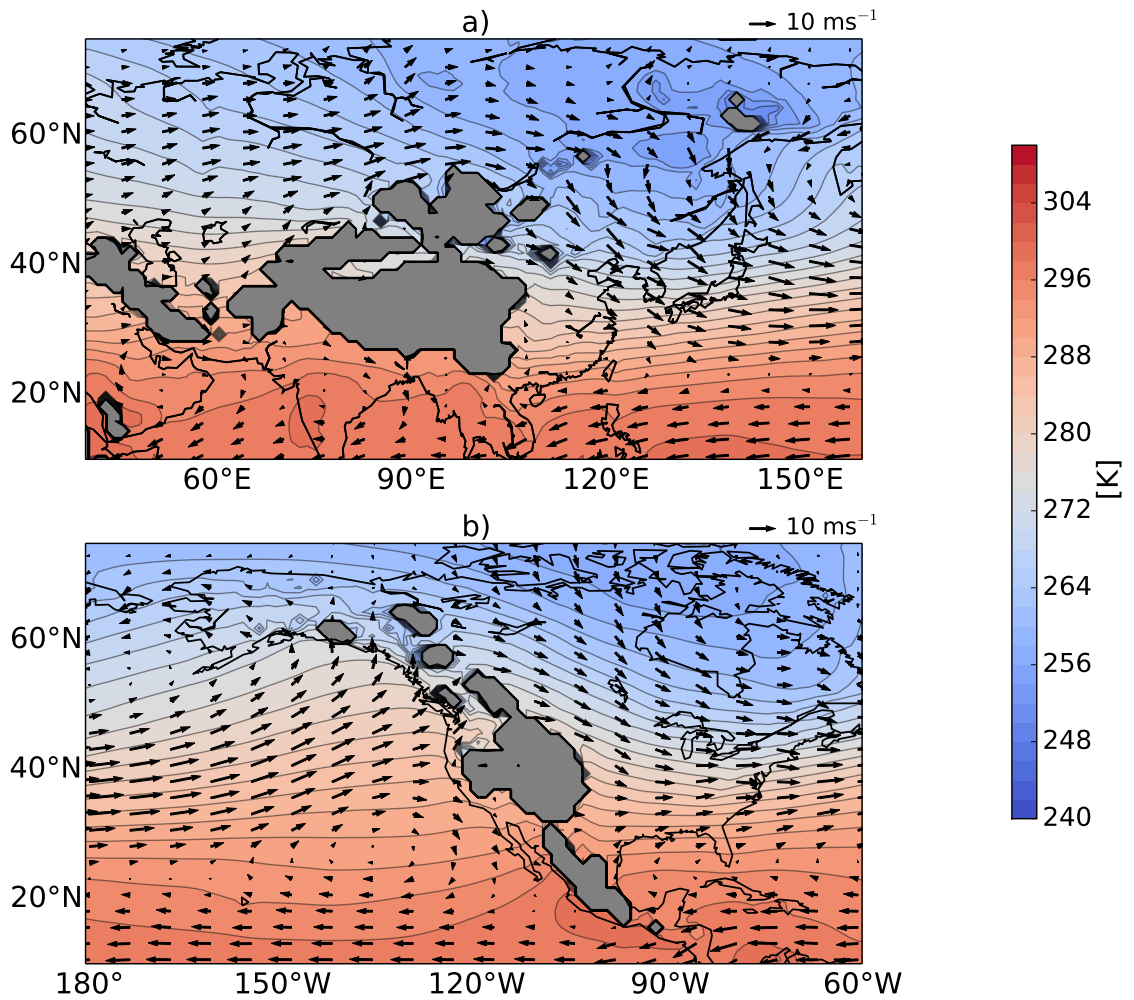
690 FIG. 6. Panels a), b) and c): advective terms in the 850hPa potential temperature variance budget from a  
 691 simulation with the moist GCM and a mountain height of 4km. Locations where topography intrudes through  
 692 850hPa are masked in gray. Panel d): the contribution of latent heating fluctuations to 850hPa potential temper-  
 693 ature variance in the same simulation.



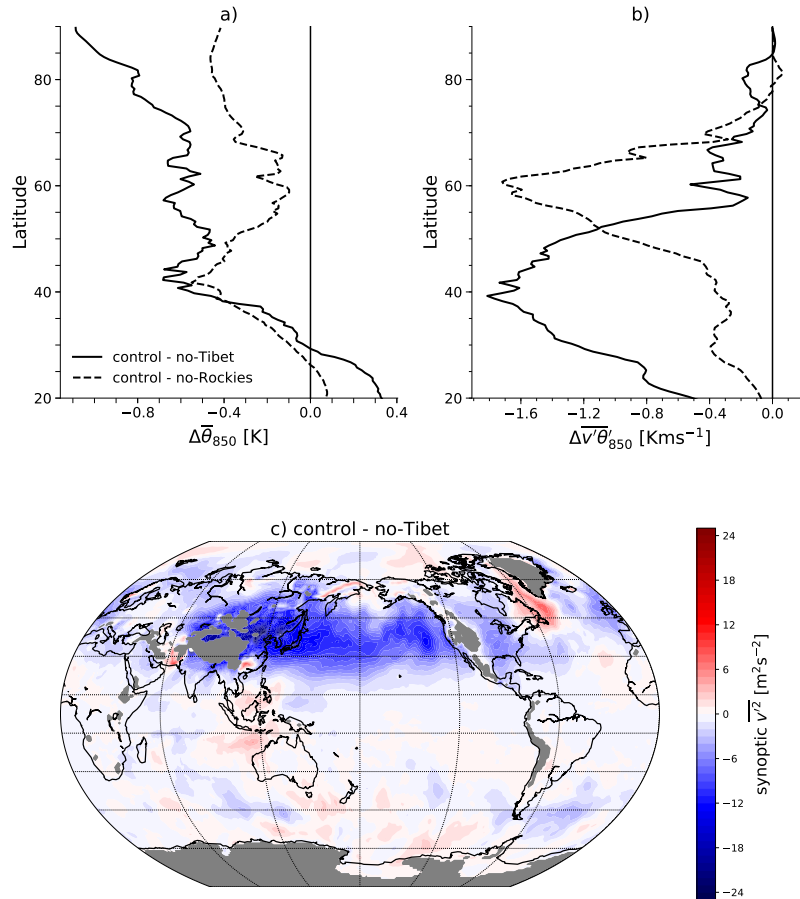
694 FIG. 7. a) Synoptic-scale variance of DJF 850hPa potential temperature in the control simulation with the  
 695 comprehensive climate model, FLOR. b) Difference in synoptic-scale variance between the control simulation  
 696 and the no-Tibet simulation. c) Difference in synoptic-scale variance between the control simulation and the  
 697 no-Rockies simulation. d) DJF squared meridional potential temperature gradients in the control simulation. e)  
 698 DJF squared meridional potential temperature gradients in the no-Tibet simulation. f) DJF squared meridional  
 699 potential temperature gradients in the no-Rockies simulation. g), h), i) DJF zonal anomalies in 850hPa potential  
 700 temperature in the same simulations. Locations where topography intrudes through 850hPa are masked in gray.



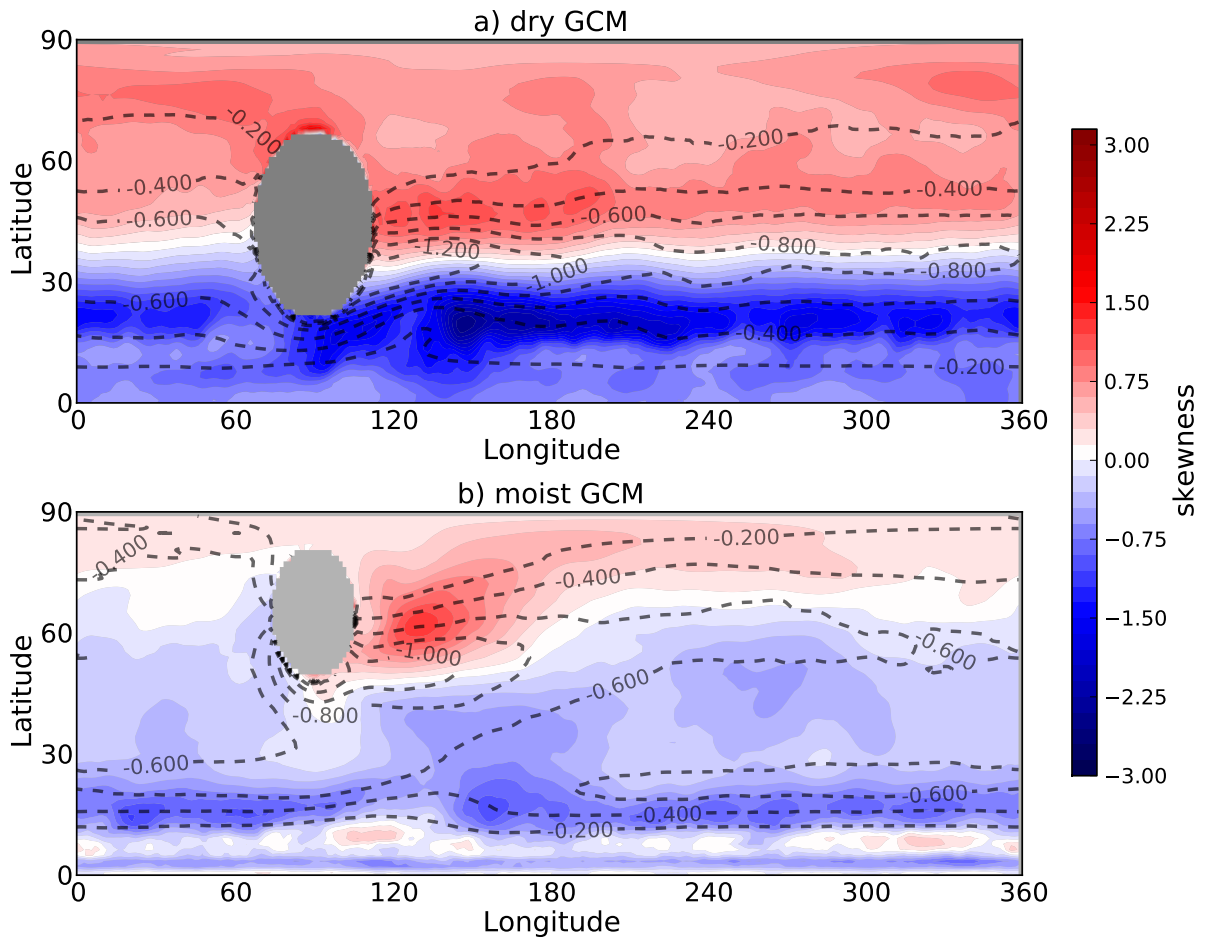
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 702 simulations with FLOR. b) Profiles taken at 50°N. Gaps in the profiles show where topography intrudes into the  
 703 850hPa level.



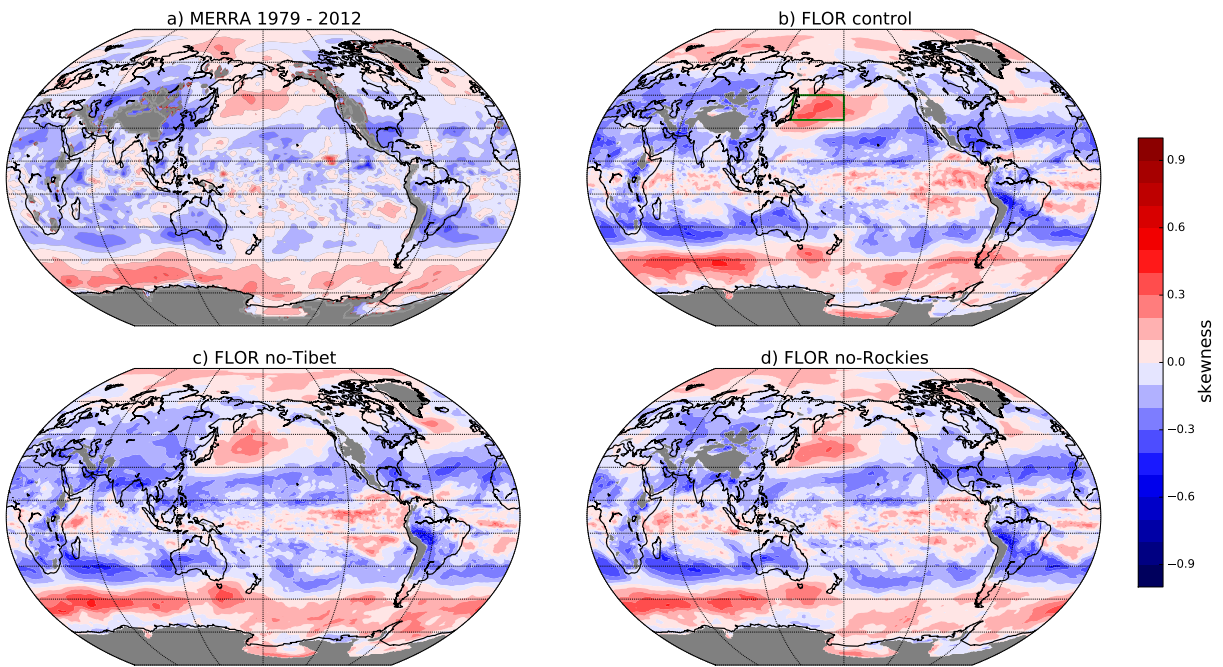
704 FIG. 9. a) DJF 850hPa temperature (contours) and total wind vectors in the vicinity of the Tibetan Plateau,  
 705 averaged over the period 1979 to 2012. Data are taken from the MERRA reanalysis dataset. Locations where  
 706 topography intrudes through 850hPa are masked in gray. b) Same for the region near the Rocky Mountains.



707 FIG. 10. a) Difference in zonal-mean  $\theta_{850}$  between the control simulation with FLOR and the no-Tibet  
 708 simulation (solid line) and difference between the control simulation and the no-Rockies simulation (dashed  
 709 line). b) Differences in transient eddy potential temperature flux in the same simulations. c) Difference in DJF  
 710 synoptic 850hPa eddy kinetic energy ( $\overline{v'^2}$ ) between the control simulation and the no-Tibet simulation.



711 FIG. 11. a) Skewness of 850hPa synoptic temperatures (colored contours) and 850hPa meridional temperature  
 712 gradients (black contours, contour interval =  $0.2\text{K}(100\text{km})^{-1}$ ) in the dry GCM simulation with  $H = 4\text{km}$ . b)  
 713 Same for the simulation with the moist GCM. The meridional gradient contour interval is  $0.2\text{K}/100\text{km}$  in both  
 714 panels.



715 FIG. 12. a) Skewness of DJF 850hPa synoptic temperatures for the period 1979-2012 in the MERRA data.  
 716 b) Skewness of DJF 850hPa synoptic temperatures in the control simulation with FLOR. c) Skewness of DJF  
 717 850hPa synoptic temperatures in the no-Tibet simulation with FLOR. d) Skewness of DJF 850hPa synoptic  
 718 temperatures in the no-Rockies simulation with FLOR.